

Imperfections – deformation and microstructures in polycrystals

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5- Elementary mechanisms

How do textures develop in a polycrystal?

- Plasticity in a single crystal
 - Slip systems
 - Schmid law
 - Critical Resolved Shear Stress (CRSS)
 - Crystal rotation
- Examples on fcc metals
- Von Mises compatibility criterion
- Other orientation mechanisms

Random questions

What is a mrd?

What is an ODF?

What is the ODF value for a random polycrystal?

What is the ODF value for a textured polycrystal?

What is the ODF value for a single-crystal?

What is the texture index? What is its value for random polycrystal, a textured polycrystal, a single-crystal?

Can you see the full texture on a pole figure?

When shall you use inverse pole figures?

5- Elementary mechanisms a- Slip systems

Slip systems

Based on experiments

- At ambient T, most of plastic deformation is accommodated through dislocation movements
- Slip always
 - Occurs on a particular set of crystallographic planes, known as *slip planes*,
 - Takes place along a consistent set of directions within these planes – these are called *slip directions*.
- The crystal structure is not altered by plastic deformation.
- Plastic deformation always occurs at constant volume.

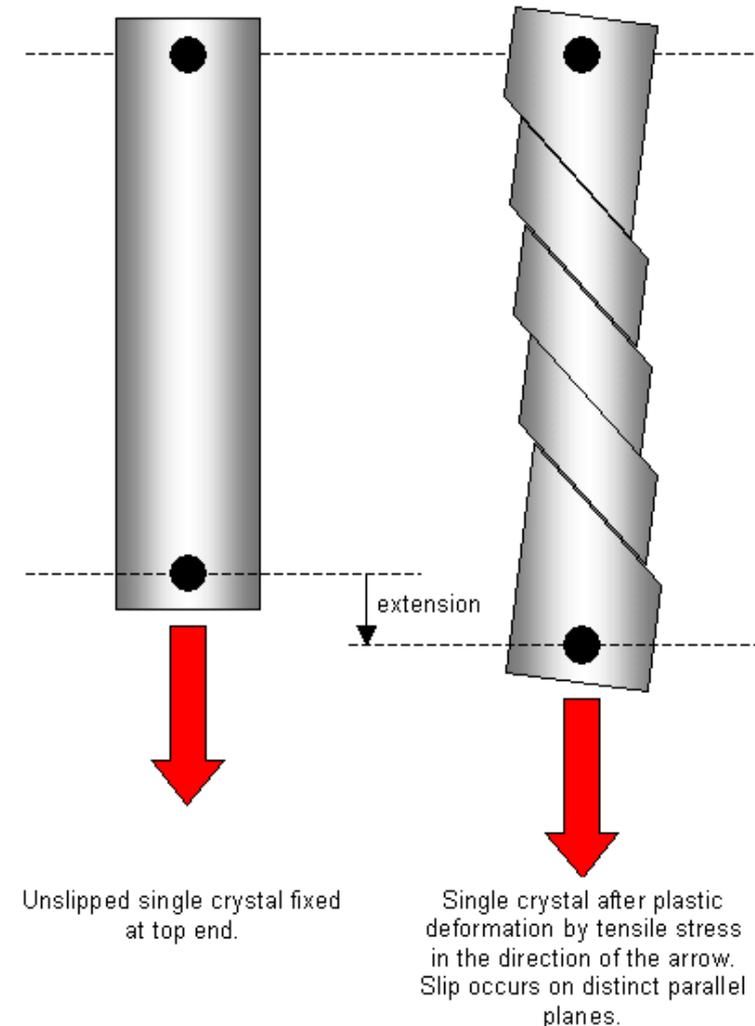
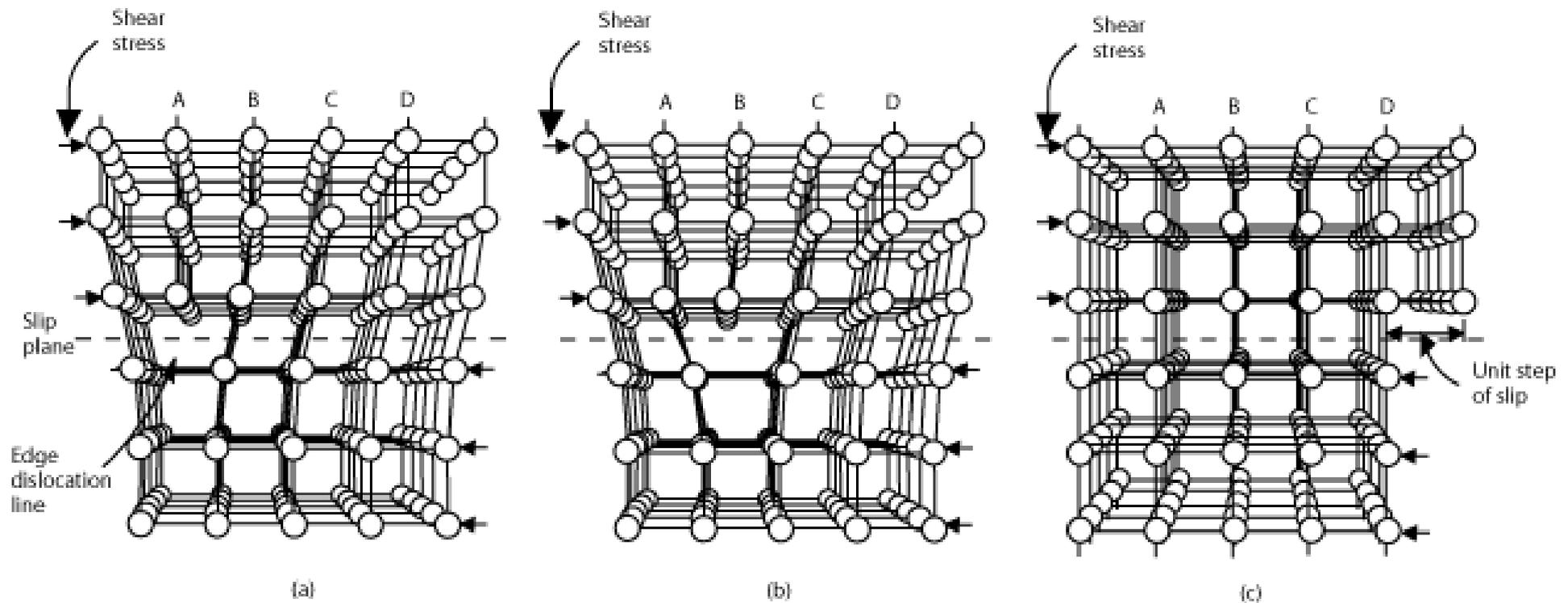


Illustration: © DoITPoMS, University of Cambridge

Effect of slip

Dislocation movement:

- One half of the crystal moves relative to the other half
- Movement in *shear*, that give rise to *deformation in shear*.



Resolved stress

- Effect of geometry on slip
- Resolved stress: applied force, projected in the direction of slip, divided by the surface of application
- Applied force, projected in the direction of slip
 $F \cos \lambda$
- Surface of application
 $A / \cos \Phi$
- Resolved stress
 $\tau = F \cos \lambda / (A / \cos \Phi)$
 $\tau = \sigma \cos \lambda \cos \Phi$
 $\tau = \sigma m ;$
- $m = \cos \lambda \cos \Phi$ is called the Schmid factor

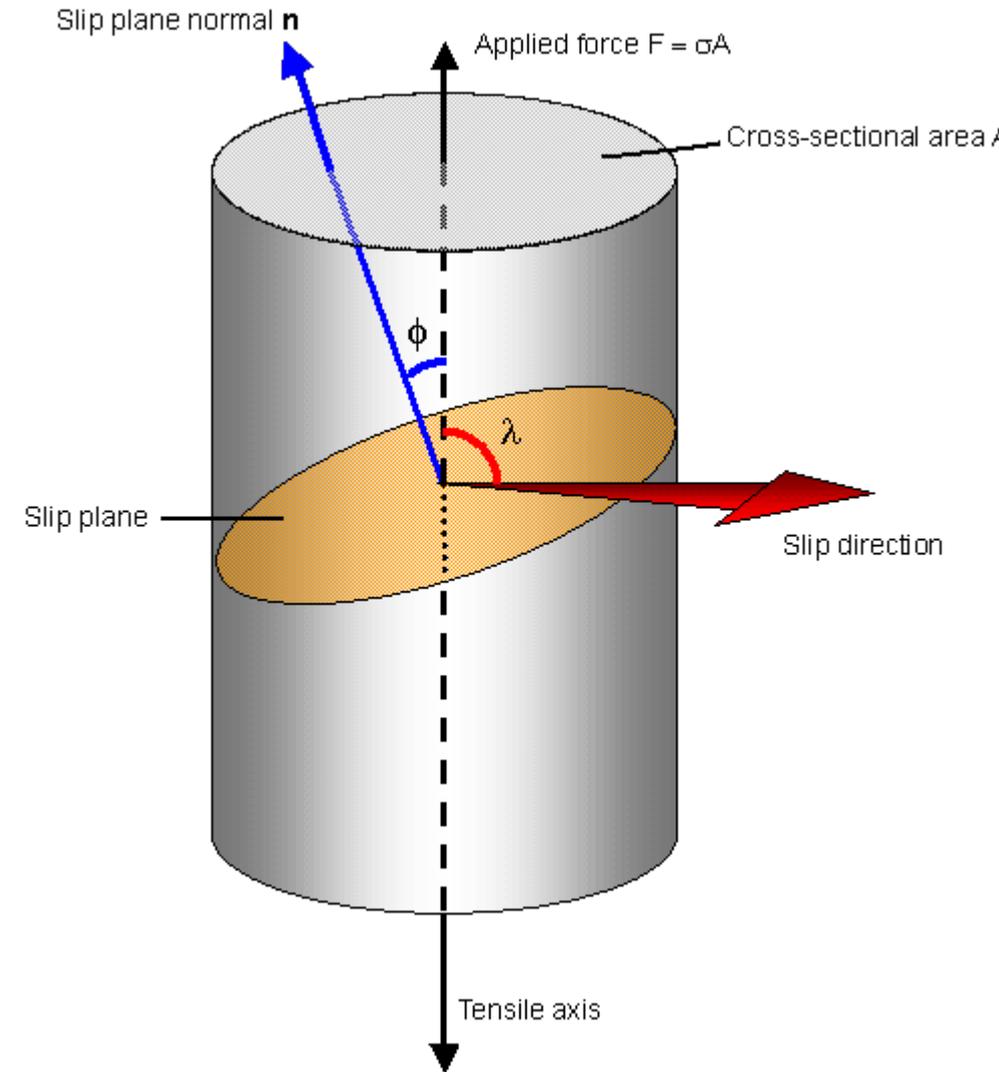


Illustration: © DoITPoMS, University of Cambridge

Schmid law:

- In a single crystal, the flow stress depends on the orientation of the crystal relative to the deformation direction
- Plastic flow begins when the resolved stress on a given slip system reaches a threshold.
- This threshold stress is independent of other stresses applied on the slip plane.
- It is called the “*Critical Resolved Shear Stress*” : the *CRSS*.

E. Schmid & W. Boas (1950),
Plasticity of Crystals, Hughes & Co., London

CRSS: illustration

Slip system with

- CRSS of 1 GPa ($\tau_c = 1$ GPa)
- $\lambda = 60^\circ$, $\Phi = 45^\circ$

Schmid factor

- $m = \cos \lambda \cos \Phi = \sqrt{2}/4 \approx 0.35$

Resolved stress

- $\tau = \sigma m$

CRSS

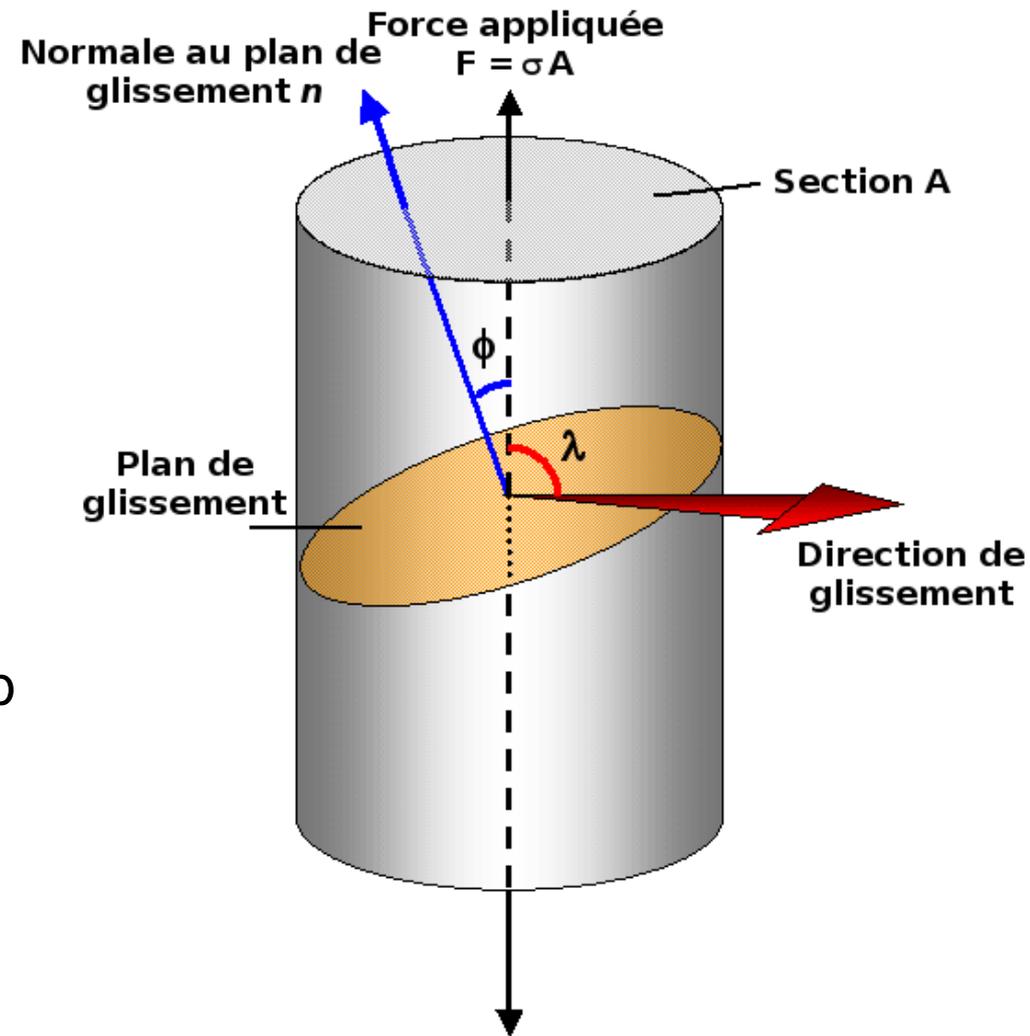
- $\tau < \tau_c$: no slip,
- $\tau = \tau_c$: slip system activated,
- $\tau > \tau_c$: should not happen, the slip system is activated when $\tau = \tau_c$.

Applied stress to get slip

- $\sigma = \tau_c / m$

Stress to be applied to have slip

- $\sigma \approx 1/0.35 \approx 2.86$ GPa



Induced rotation: tension

The slip direction rotates towards the tension direction.

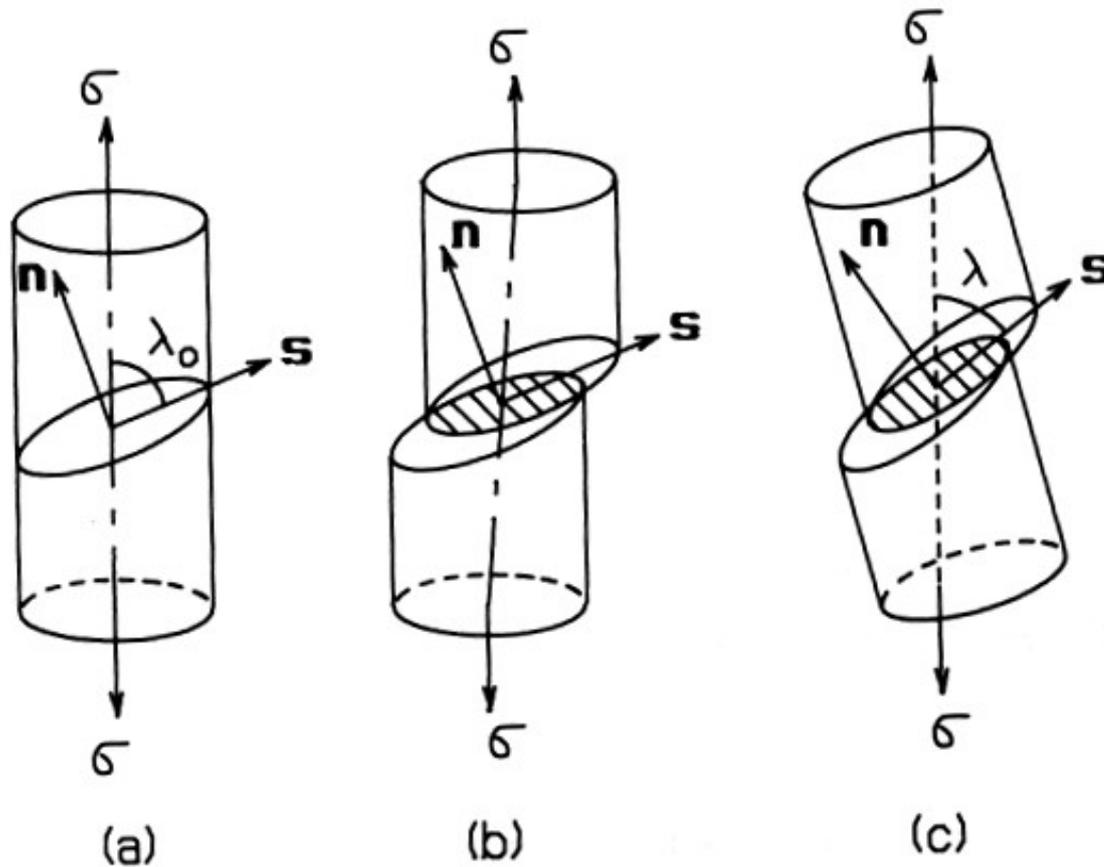


Fig. 10.5 Rotation of crystal lattice in tension

Induced rotation: compression

The slip plane becomes orthogonal to the compression direction.

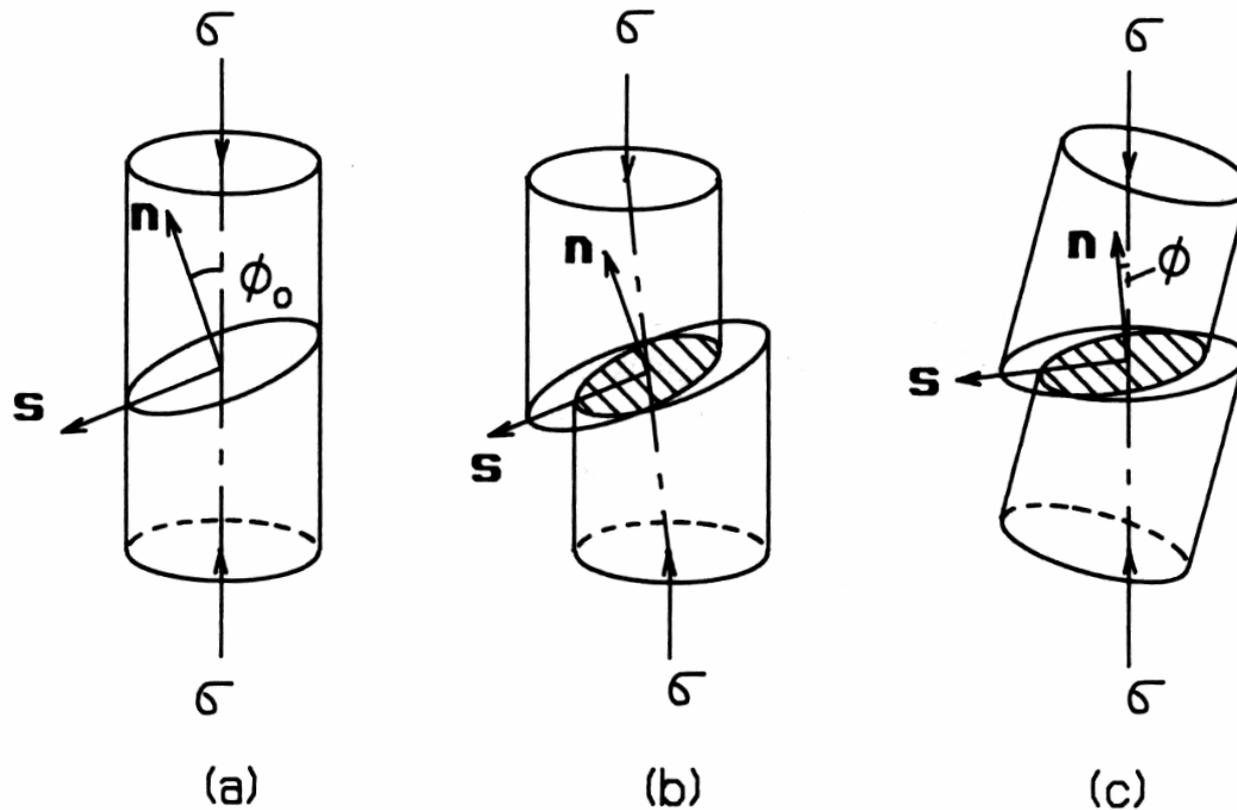
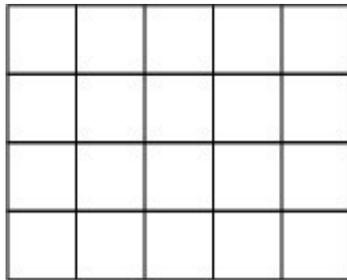
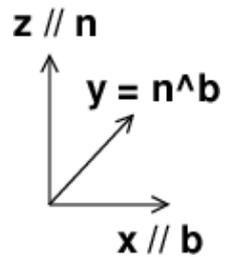
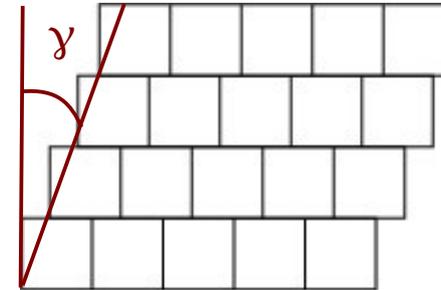


Fig. 10.6 Rotation of crystal lattice in compression

Equations



Slip // \mathbf{b} , in the plane normal to \mathbf{n}



Deformation

$$\vec{du}(x, y, z) = \begin{cases} \gamma z \\ 0 \\ 0 \end{cases}$$

Deformation gradient

$$F_{ij} = \frac{\partial du_i}{\partial x_j}$$

$$F = \begin{bmatrix} 0 & 0 & \gamma \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Kinematic analysis

$$F = \epsilon + \omega$$

$$\epsilon = \begin{bmatrix} 0 & 0 & \gamma/2 \\ 0 & 0 & 0 \\ \gamma/2 & 0 & 0 \end{bmatrix}$$

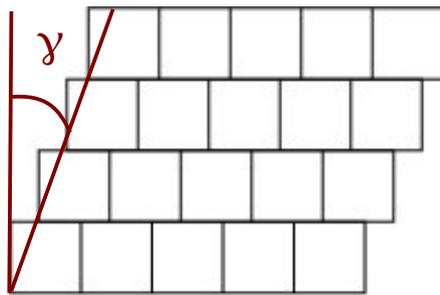
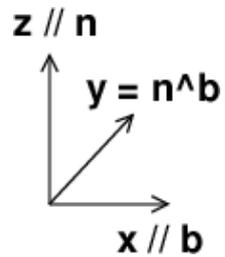
Shear

- Trace = 0
- No volume change

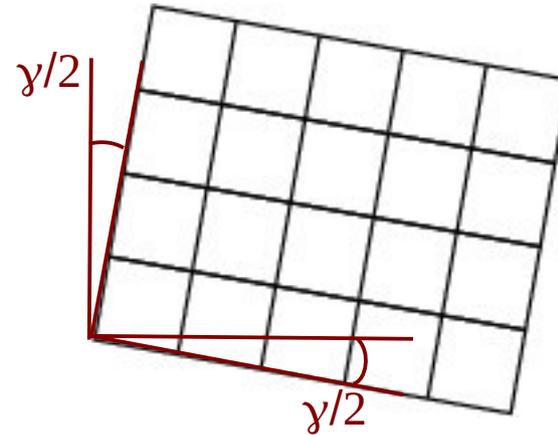
$$\omega = \begin{bmatrix} 0 & 0 & \gamma/2 \\ 0 & 0 & 0 \\ -\gamma/2 & 0 & 0 \end{bmatrix}$$

Rotation due to slip

Geometric interpretation

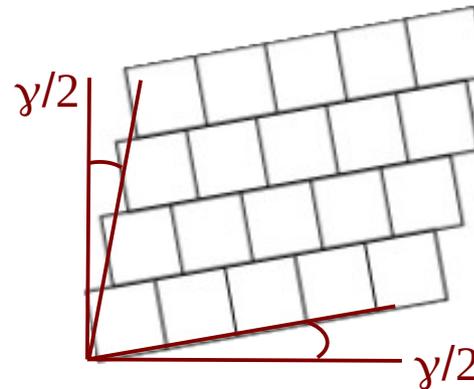


$$F = \begin{bmatrix} 0 & 0 & \gamma \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$



$$\omega = \begin{bmatrix} 0 & 0 & \gamma/2 \\ 0 & 0 & 0 \\ -\gamma/2 & 0 & 0 \end{bmatrix}$$

= +



$$\epsilon = \begin{bmatrix} 0 & 0 & \gamma/2 \\ 0 & 0 & 0 \\ \gamma/2 & 0 & 0 \end{bmatrix}$$

General setting

- Slip system with, \mathbf{b} , burgers vector or slip direction, \mathbf{n} normal to slip plane
- Homogeneous deformation (many dislocations uniformly distributed in the slip plane)
- Total macroscopique deformation: γ

Movement

$$\mathbf{u}(x, y, z) = \gamma (\mathbf{r} \cdot \mathbf{n}) \mathbf{l}$$

$$\mathbf{l} = \mathbf{b} / \|\mathbf{b}\|$$

$$\mathbf{r} = \begin{bmatrix} x \\ y \\ z \end{bmatrix}$$

Infinitesimal deformation

$$\epsilon_{ij} = \frac{1}{2} \gamma (n_j l_i + n_i l_j)$$

Infinitesimal rotation

$$\omega_{ij} = \frac{1}{2} \gamma (n_j l_i - n_i l_j)$$

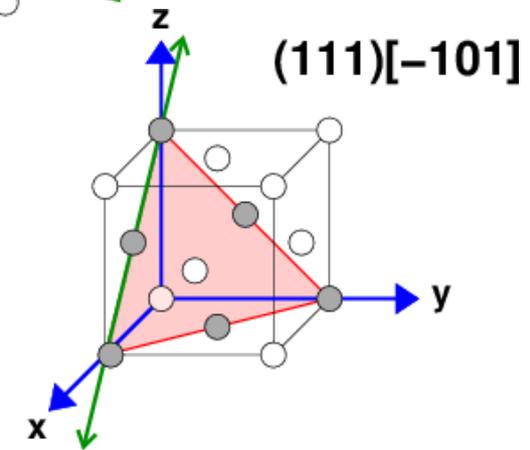
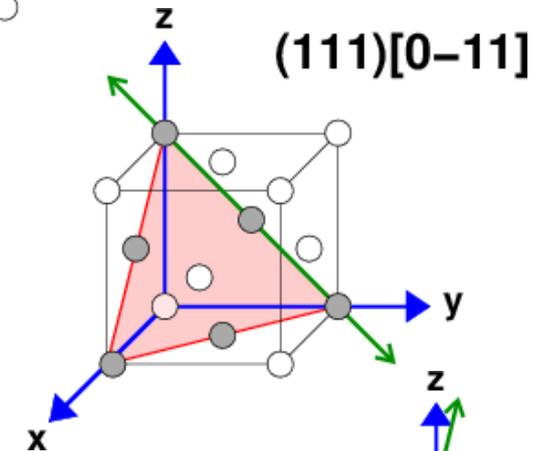
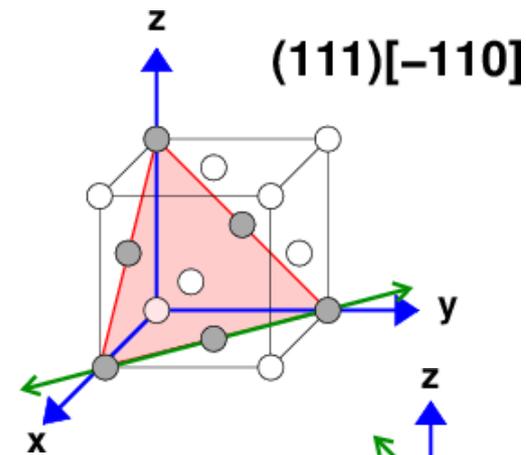
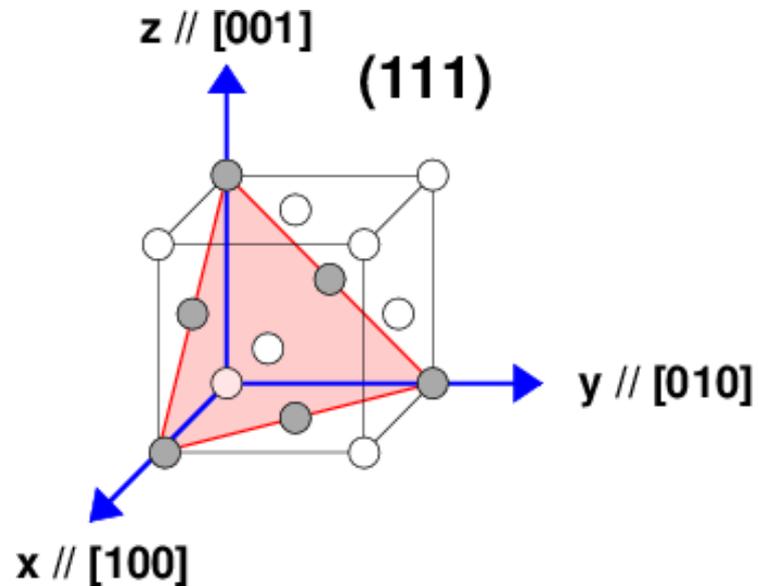
5- Elementary mechanisms *b*- Face-centered-cubic metals

Fcc - slip systems

Slip systems in the fcc structure:

- Slip direction: $\langle 110 \rangle$
- Slip plane: $\{111\}$

Slip system: $(111)[-110]$



Fcc - slip systems

- 4 {111} planes: (111), ($\bar{1}11$), ($1\bar{1}1$), ($11\bar{1}$) ;
- For each: 3 directions $\langle 110 \rangle$;
- 12 slip systems, each can operate in 2 directions (+ ou -).

Number	a1	a2	a3	b1	b2	b3
Plane	(111)	(111)	(111)	($\bar{1}\bar{1}1$)	($\bar{1}\bar{1}1$)	($\bar{1}\bar{1}1$)
Direction	[$0\bar{1}1$]	[$10\bar{1}$]	[$\bar{1}10$]	[011]	[101]	[$1\bar{1}0$]

Number	c1	c2	c3	d1	d2	d3
Plane	($\bar{1}11$)	($\bar{1}11$)	($\bar{1}11$)	($1\bar{1}1$)	($1\bar{1}1$)	($1\bar{1}1$)
Direction	[$0\bar{1}1$]	[$\bar{1}0\bar{1}$]	[110]	[011]	[$10\bar{1}$]	[$\bar{1}\bar{1}0$]

Let's consider a fcc crystal with stress applied parallel to [001]:

- Evaluate the Schmid factor for all 12 slip systems,
- How many will be activated?

Number	a1	a2	a3	b1	b2	b3
Plane	(111)	(111)	(111)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)
Direction	[0 $\bar{1}$ 1]	[10 $\bar{1}$]	[$\bar{1}$ 10]	[011]	[101]	[1 $\bar{1}$ 0]

a1 system

- Stress direction: $\mathbf{s} = (0;0;1)$
- Slip direction: $\mathbf{l} = 1/\sqrt{2} (0;-1;1)$
- Normal to slip plane: $\mathbf{n} = 1/\sqrt{3} (1;1;1)$
- Schmid factor: $m = \cos \lambda \cos \Phi$
- $m = \mathbf{s} \cdot \mathbf{l} \mathbf{s} \cdot \mathbf{n}$
- $m = 1/\sqrt{6}$

Practice 1

Let's consider a fcc crystal with stress applied parallel to [001]:

- Evaluate the Schmid factor for all 12 slip systems,
- How many will be activated?

Number	a1	a2	a3	b1	b2	b3
Plane	(111)	(111)	(111)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)
Direction	[0 $\bar{1}\bar{1}$]	[10 $\bar{1}$]	[$\bar{1}\bar{1}$ 0]	[011]	[101]	[1 $\bar{1}$ 0]
Schmid F.	$1/\sqrt{6}$	$-1/\sqrt{6}$	0	$1/\sqrt{6}$	$1/\sqrt{6}$	0

Number	c1	c2	c3	d1	d2	d3
Plane	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)
Direction	[0 $\bar{1}\bar{1}$]	[$\bar{1}$ 0 $\bar{1}$]	[110]	[011]	[10 $\bar{1}$]	[$\bar{1}\bar{1}$ 0]
Schmid F.	$1/\sqrt{6}$	$-1/\sqrt{6}$	0	$1/\sqrt{6}$	$-1/\sqrt{6}$	0

8
systems
activated

Practice 2

Let's consider a fcc crystal with stress applied parallel to $[011]$:

- Same questions...

Number	a1	a2	a3	b1	b2	b3
Plane	(111)	(111)	(111)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)
Direction	$[0\bar{1}1]$	$[10\bar{1}]$	$[\bar{1}10]$	$[011]$	$[101]$	$[1\bar{1}0]$

Number	c1	c2	c3	d1	d2	d3
Plane	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($1\bar{1}\bar{1}$)	($1\bar{1}\bar{1}$)	($1\bar{1}\bar{1}$)
Direction	$[0\bar{1}1]$	$[\bar{1}0\bar{1}]$	$[110]$	$[011]$	$[10\bar{1}]$	$[\bar{1}\bar{1}0]$

Practice 2

Let's consider a fcc crystal with stress applied parallel to [011]:

- Same questions...

Number	a1	a2	a3	b1	b2	b3
Plane	(111)	(111)	(111)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)
Direction	[0 $\bar{1}$ 1]	[10 $\bar{1}$]	[$\bar{1}$ 10]	[011]	[101]	[1 $\bar{1}$ 0]
Schmid F.	0	-1/ $\sqrt{6}$	1/ $\sqrt{6}$	0	0	0

Number	c1	c2	c3	d1	d2	d3
Plane	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)	(1 $\bar{1}\bar{1}$)
Direction	[0 $\bar{1}$ 1]	[$\bar{1}$ 0 $\bar{1}$]	[110]	[011]	[10 $\bar{1}$]	[$\bar{1}$ 10]
Schmid F.	0	-1/ $\sqrt{6}$	1/ $\sqrt{6}$	0	0	0

4
systems
activated

Choice of slip system for fcc metals

Graphical representation

Active slip system in fcc,
as a function of the
direction of applied stress

- signs above slip system numbers:
 - opposite direction relative to that of the table
 - for instance $[-1-10]$ instead of $[110]$.

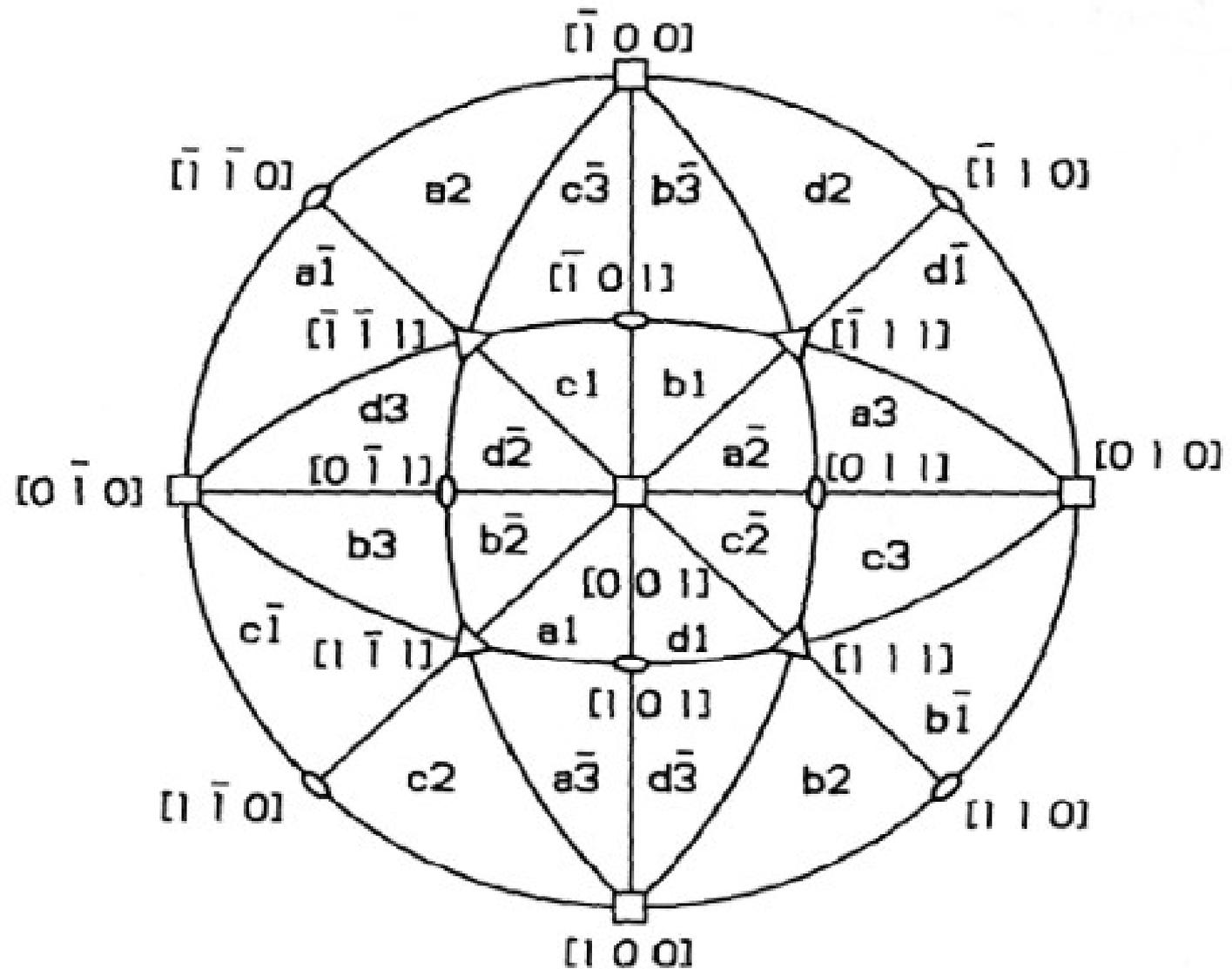


Image A. Rollett,
after Kahn

More complex cases

For fcc metals, there is only one type of slip systems → the CRSS is not needed to evaluate with slip plane and directions are activated first.

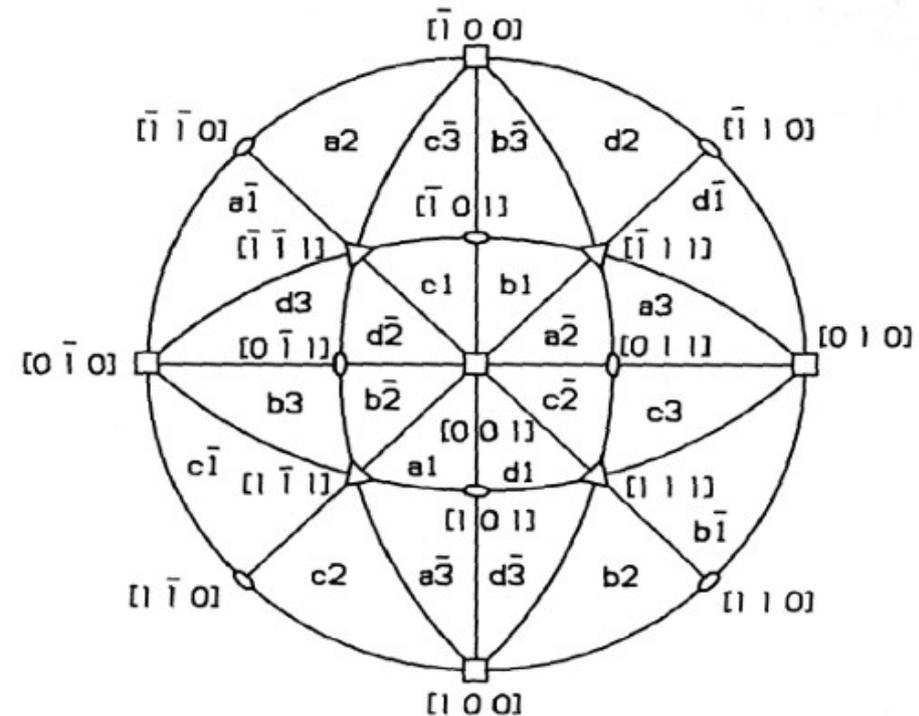
In other systems, you may have different types of slip systems, with different CRSS. Then, you have to look at the resolved shear stress on all slip systems and see which one is closest to the CRSS.

Sometimes, “secondary slip systems” can be activated. Their CRSS is higher, the resolved shear stress is higher, but these systems are required for plastic deformation to work → *von Mises compatibility criterion* (later in the course)

Consider a fcc crystal with applied stress parallel to $[102]$:

- Identify the activated slip systems,
- For each, establish the expression of the rotation tensor,
- Assume that deformation is evenly distributed between all active slip systems and evaluate to complete rotation tensor,
- Describe the effect of the rotation on the crystal

a1	a2	a3	b1	b2	b3
(111)	(111)	(111)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)
$[0\bar{1}1]$	$[10\bar{1}]$	$[\bar{1}10]$	$[011]$	$[101]$	$[1\bar{1}0]$
c1	c2	c3	d1	d2	d3
($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	($\bar{1}\bar{1}\bar{1}$)	(111)	(111)	(111)
$[0\bar{1}1]$	$[\bar{1}0\bar{1}]$	$[110]$	$[011]$	$[10\bar{1}]$	$[\bar{1}\bar{1}0]$



- Activated systems : a1 et d1

$$(111)[0\bar{1}1]$$

$$\mathbf{n} = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ 1 \\ 1 \end{bmatrix} \quad \mathbf{l} = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$$

$$\omega_1 = \frac{\gamma}{2\sqrt{6}} \begin{bmatrix} 0 & 1 & -1 \\ -1 & 0 & -2 \\ 1 & 2 & 0 \end{bmatrix}$$

$$(1\bar{1}1)[011]$$

$$\mathbf{n} = \frac{1}{\sqrt{3}} \begin{bmatrix} 1 \\ -1 \\ 1 \end{bmatrix} \quad \mathbf{l} = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 1 \\ 1 \end{bmatrix}$$

$$\omega_2 = \frac{\gamma}{2\sqrt{6}} \begin{bmatrix} 0 & -1 & -1 \\ 1 & 0 & 2 \\ 1 & -2 & 0 \end{bmatrix}$$

$$\omega_{tot} = \frac{\gamma}{2\sqrt{6}} \begin{bmatrix} 0 & 0 & -1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{bmatrix}$$

ω_{ij} = anti-symmetric component of the deformation gradient
= infinitesimal rotation tensor

Effect on a vector

$$du_i = \omega_{ij} u_j$$

$$\begin{bmatrix} du_x \\ du_y \\ du_z \end{bmatrix} = \begin{bmatrix} 0 & -\omega_z & +\omega_y \\ +\omega_z & 0 & -\omega_x \\ -\omega_y & +\omega_x & 0 \end{bmatrix} \begin{bmatrix} u_x \\ u_y \\ u_z \end{bmatrix}$$

\mathbf{u} is the vector and $d\mathbf{u}$ the displacement. ω is not a “true” rotation matrix, but a tensor.

5- Elementary mechanisms c- Von Mises compatibility criterion

Von Mises compatibility criterion

Von Mises compatibility criterion

- All strains must be accommodated by plastic deformation,
- All component of the strain tensor must be accommodated,
- Plasticity = constant volume deformation
- $\epsilon_{11} + \epsilon_{22} + \epsilon_{33} = 0$
- 5 independent component left : $\epsilon_{11}, \epsilon_{22}, \epsilon_{12}, \epsilon_{13}, \epsilon_{23}$
- The material needs at least 5 independent deformation mechanisms

$$\epsilon_{ij}^T = \sum_k \frac{1}{2} \gamma^k (n_j^k l_i^k + n_i^k l_j^k)$$

Otherwise:

- There must be another mechanism coming in, otherwise the crystal will break.

MgO is a cubic material with a B1 structure, like NaCl

- The $\{110\}\langle\bar{1}\bar{1}0\rangle$ slip system is dominant, and has six symmetry equivalents
 - 1 : $(110)[\bar{1}\bar{1}0]$
 - 2 : $(1\bar{1}0)[110]$
 - 3 : $(011)[0\bar{1}\bar{1}]$
 - 4 : $(0\bar{1}\bar{1})[011]$
 - 5 : $(101)[\bar{1}0\bar{1}]$
 - 6 : $(10\bar{1})[101]$
- Are they sufficient to accommodate any form of deformation?

$$\epsilon_{ij} = \frac{1}{2}\gamma (n_j l_i + n_i l_j)$$

$$\omega_{ij} = \frac{1}{2}\gamma (n_j l_i - n_i l_j)$$

MgO is a cubic material with a B1 structure, like NaCl

- The $\{110\}\langle 110\rangle$ slip system is dominant, and has six symmetry equivalents.
- Are they sufficient to accommodate any form of deformation?

$$(110) [\bar{1}10]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} -1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$(011) [0\bar{1}1]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$(101) [\bar{1}01]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} -1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

No component
of strain off-
diagonal

$$(\bar{1}10) [110]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

$$(01\bar{1}) [011]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

$$(10\bar{1}) [101]$$

$$\epsilon = \frac{\gamma}{2} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

Need additional
systems such as
 $\{100\}\langle 1\bar{1}0\rangle$

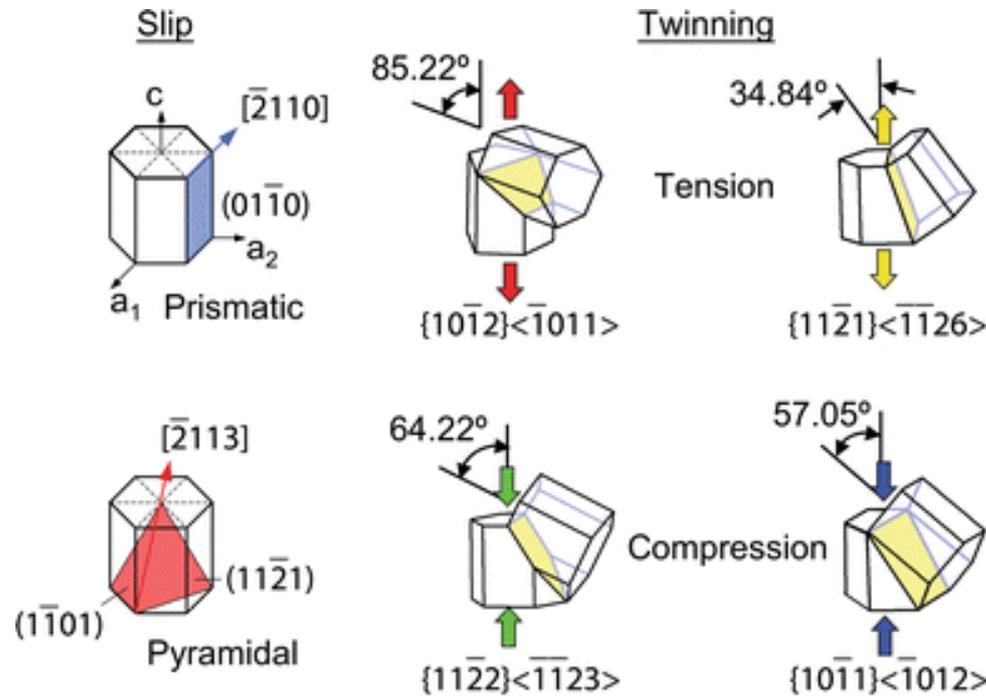
5- Elementary mechanisms c- Other mechanisms

Other mechanisms

Dislocations are not the only mechanism that can rotate grains and generate grain orientations.

Other important mechanism: twinning.

Twin: two separate crystals share some of the same crystal lattice points in a symmetrical manner.



Example: deformation mechanisms and orientation in zirconium.

Padilla et al, *Metallurgical and Materials Transactions A*, 2007

